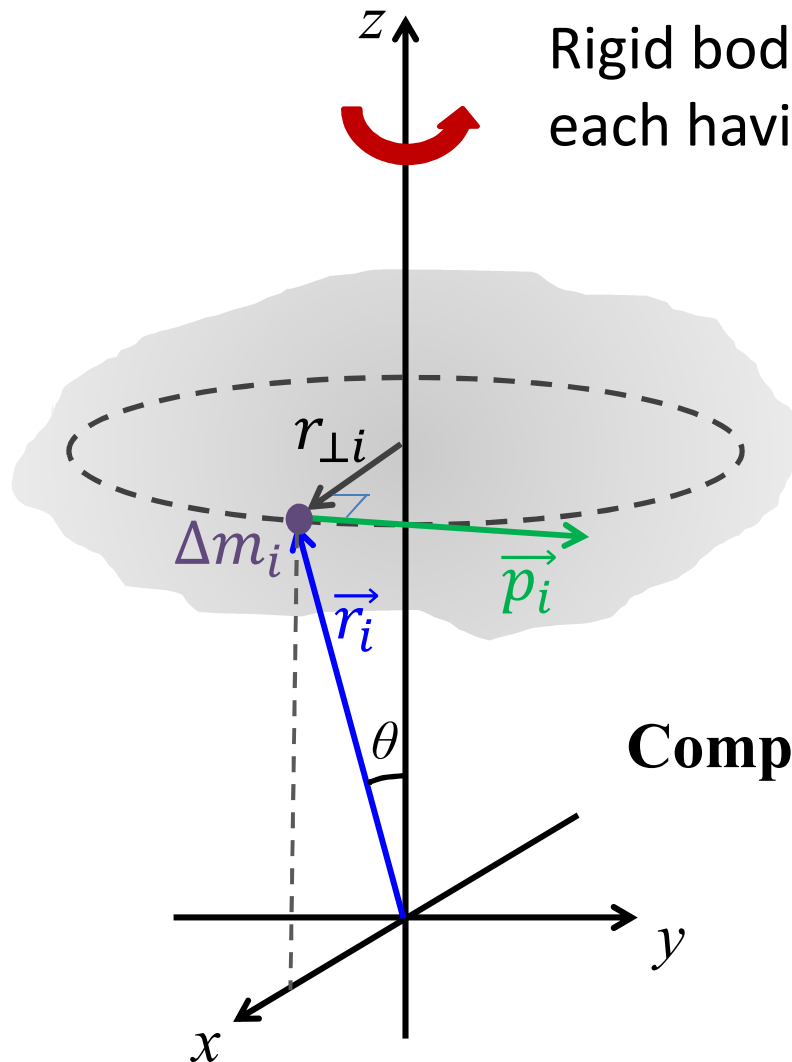


CHAPTER 11 (TORQUE, ANGULAR MOMENTUM)

1. Angular Momentum \vec{L}
 - a. Rigid body
 - b. Conservation of \vec{L}
2. Applications
 - a. Gyroscope

ANGULAR MOMENTUM OF A RIGID BODY (ARBITRARY SHAPE AND AXIS POSITION)



Rigid body taken as a system of many elements, each having mass Δm .

About point O:
$$\vec{L}_O = \sum_i \vec{l}_i$$

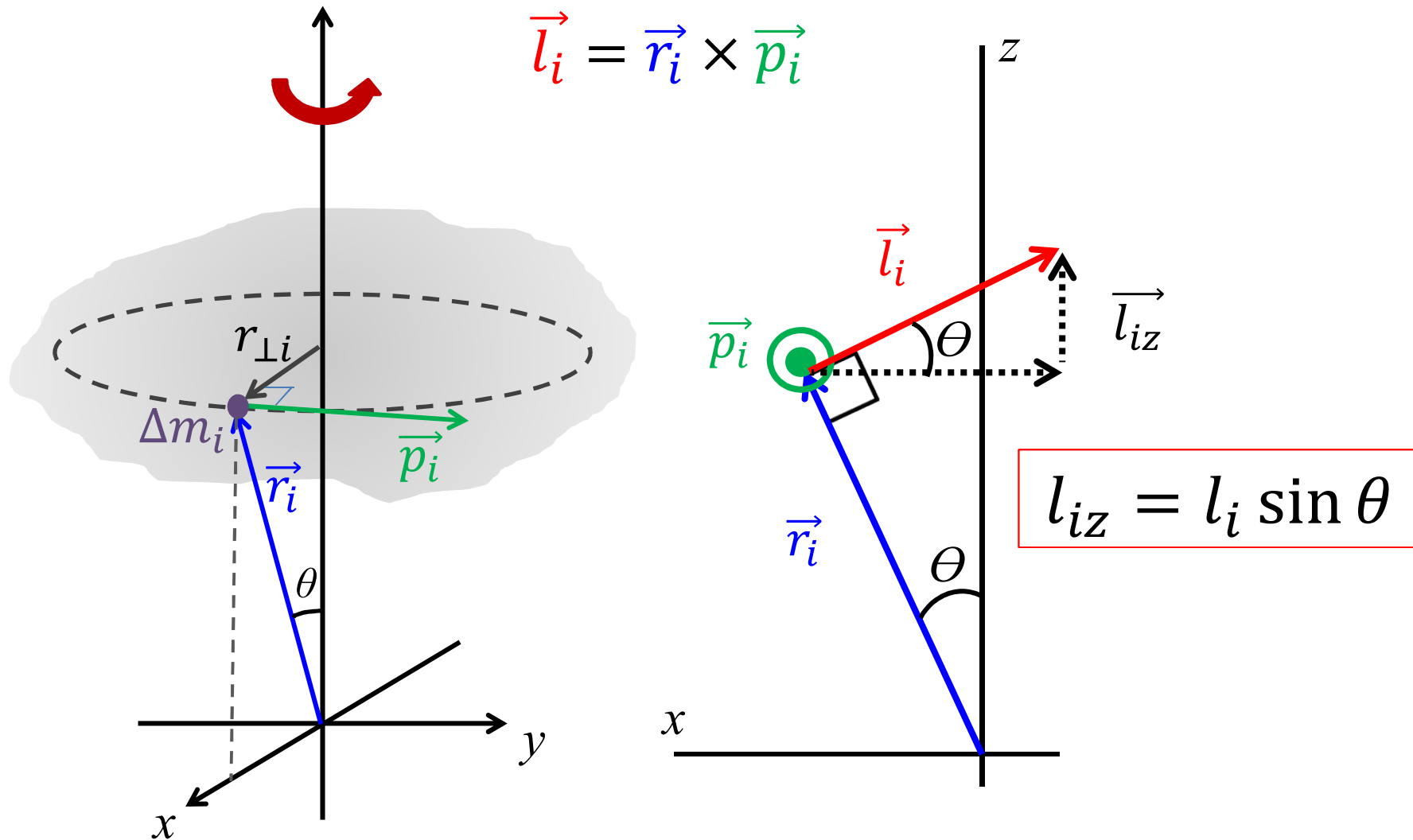
$$\vec{l}_i = \vec{r}_i \times \vec{p}_i$$

$$\vec{r}_i \perp \vec{p}_i \Rightarrow |\vec{l}_i| = r_i \Delta m v_i$$

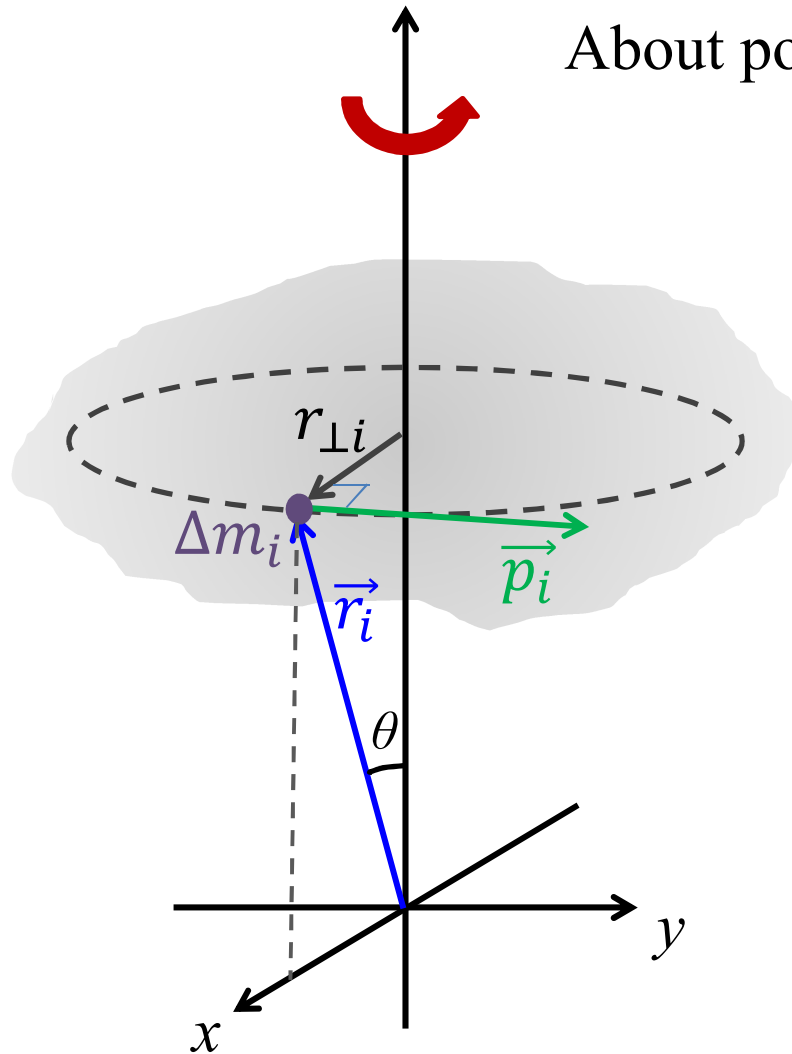
Component L_{Oz} of \vec{L}_O along rotation (z) axis?

$$L_{Oz} = \sum_i l_{iz}$$

COMPONENT OF l_i ALONG ROTATION AXIS Z



ANGULAR MOMENTUM OF A RIGID BODY - COMPONENT ALONG ROTATION AXIS -



About point O: $\vec{L}_O = \sum_i \vec{l}_i$ $|\vec{l}_i| = r_i \Delta m v_i$

$$l_{iz} = l_i \sin \theta$$

$$= r_i \sin \theta \Delta m v_i = r_{\perp i} \Delta m v_i$$

But $v_i = r_{\perp i} \omega$

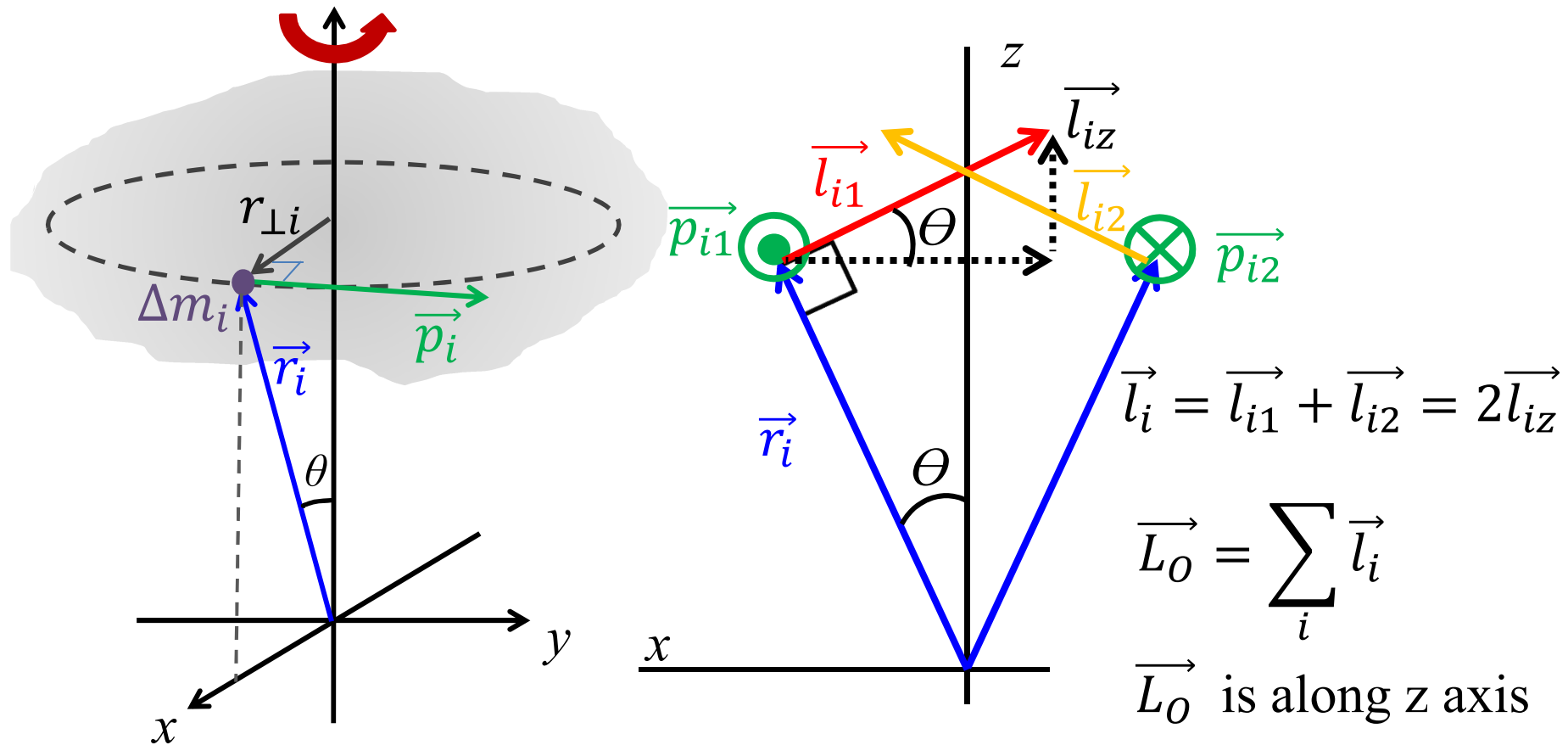
Then, $l_{iz} = r_{\perp i} \Delta m r_{\perp i} \omega = \Delta m r_{\perp i}^2 \omega$

$$L_{Oz} = \sum_i l_{iz} = \omega \sum_i \Delta m r_{\perp i}^2$$

$$\Rightarrow L_z = I_z \omega$$

Characteristic of axis z, not of point O

ANGULAR MOMENTUM OF A SYMMETRIC RIGID BODY ROTATING ABOUT AN AXIS OF SYMMETRY



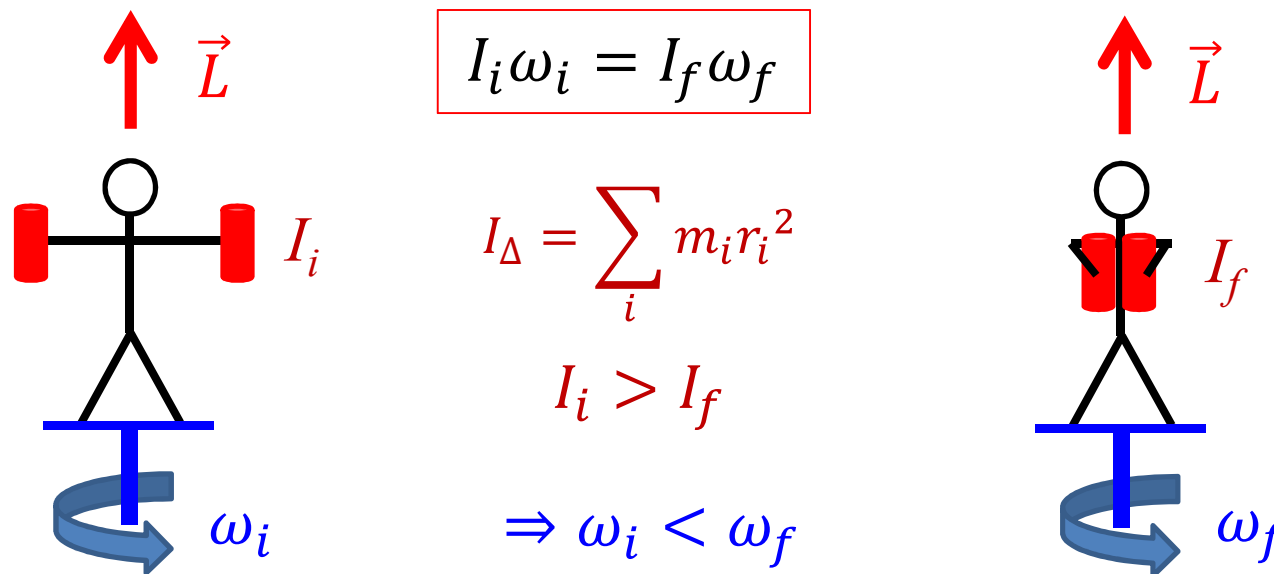
If the axis of rotation is also **an axis of symmetry** of the body:

$$\vec{L} = I_z \vec{\omega}$$

\vec{L} is along z axis; characteristic of the axis

CONSERVATION OF ANGULAR MOMENTUM

If the component of the net *external* torque on a system along a certain axis is zero, then the component of the angular momentum of the system along that axis cannot change, no matter what changes take place within the system.



Initial State (i): The man has a relatively large rotational inertia about the rotation axis and a relatively small angular speed.

Final State (f): By decreasing his rotational inertia, the man increases his angular speed. The angular momentum of the rotating system remains unchanged.

FORMULA SUMMARY

Torque: $\vec{\tau} = \vec{r} \times \vec{F}$

Angular momentum: $\vec{l} = \vec{r} \times \vec{p}$

Newton's 2nd law: $\vec{\tau}_{net} = \frac{d\vec{L}}{dt}$ Pure rotation: $\tau_{net} = I \frac{d\omega}{dt} = I\alpha$

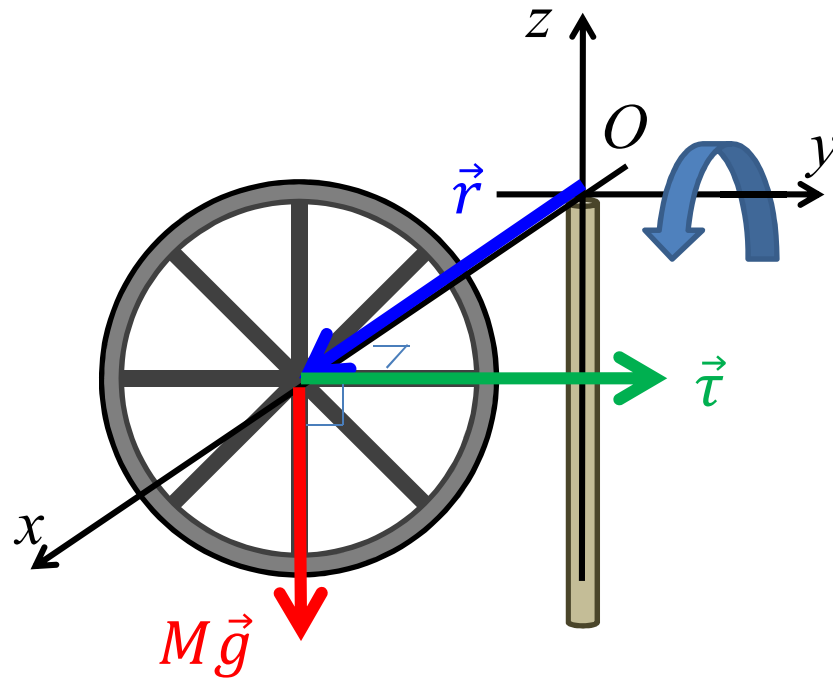
Rigid body, rotation about fixed axis z: $L_z = I_z \omega$

Symmetric body rotation, about symmetry axis z: $\vec{L} = I_z \vec{\omega}$

General motion (translation + rotation)

→ use Newton's 2nd law in **both** linear and angular forms

GYROSCOPE NOT SPINNING



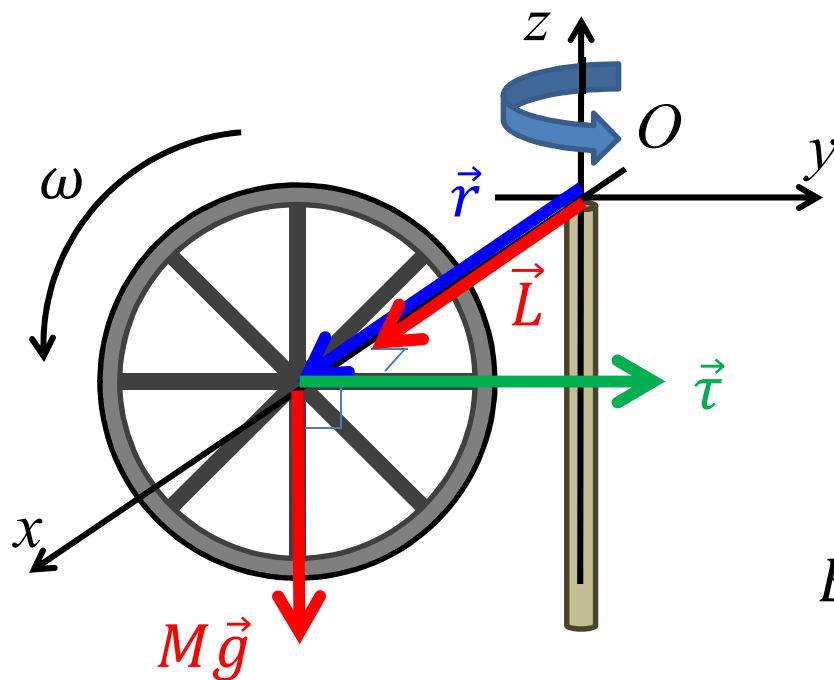
$$\vec{\tau} = \frac{d\vec{L}}{dt}$$

$$|\vec{\tau}| = MGr \sin \pi/2 = Mgr$$

A non-spinning gyroscope falls by rotating in an xz plane because of torque τ .

PRECESSION OF A GYROSCOPE

A rapidly spinning gyroscope, with angular momentum \vec{L} , precesses around the z axis. Its precessional motion is in the xy plane.



Symmetric body rotating about symmetry axis Ox : $\vec{L}(t) = I\vec{\omega}(t)$

$$\Rightarrow \frac{d\vec{L}}{dt} = I \frac{d\vec{\omega}}{dt} = \overrightarrow{\tau_{gravity}}$$

But: $|\vec{L}(t)| = I|\vec{\omega}(t)| = \text{Constant!}$

The only way the vector $\vec{L}(t)$ can change with time is by changing its orientation! How does it change?

PRECESSION RATE OF A GYROSCOPE

$d\vec{L} = \vec{\tau} dt \rightarrow \vec{L}$ Changes in the direction of $\vec{\tau}$

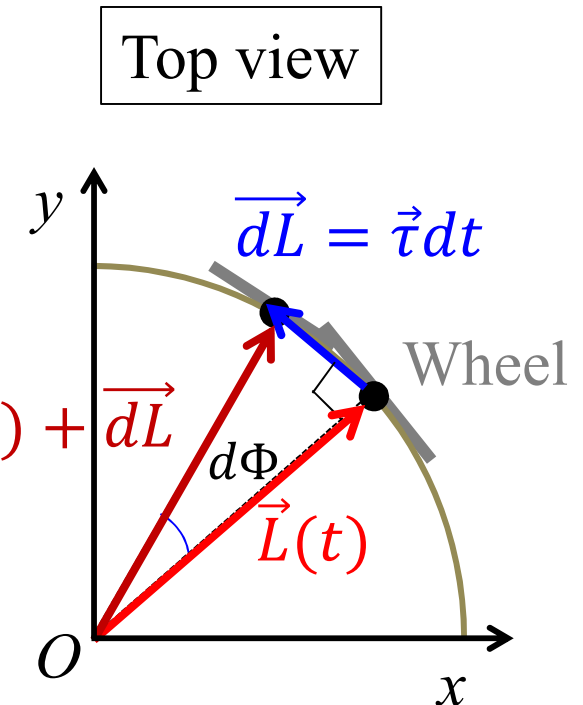
$$|d\vec{L}| = |\vec{L}| d\phi$$

$$\Rightarrow \frac{d\phi}{dt} = \frac{1}{|\vec{L}|} \frac{|d\vec{L}|}{dt} = \frac{|\vec{\tau}|}{|\vec{L}|} = \frac{Mgr}{I\omega}$$

$$\Omega = \frac{Mgr}{I\omega}$$

The change of \vec{L} , which is induced by the torque $\vec{\tau}$ leads to a rotation of \vec{L} about O , with angular velocity Ω (Precession rate).

$$\vec{L}(t + \Delta t) = \vec{L}(t) + d\vec{L}$$



PRECESSION OF A GYROSCOPE

Note: precession induces a new component of the angular momentum: $\vec{l}_{prec} = I_z \vec{\Omega}$ along the z axis! But negligible if ω is large (rapid spinning) : $|\vec{l}_{prec}| \ll |\vec{L}|$.

